Initial Boundary Value Problem for a Generalized Zakharov Equations

Shujun You, Xiaoqi Ning

Abstract—This paper considers the existence of the generalized solution to the initial boundary value problem for a class of generalized Zakharov equation in (2 + 1) dimensions. By a priori integral estimates and the Galerkin method, the author establishes global in time existence of the solution to the problem.

Index Terms—generalized solution, generalized Zakharov equations, initial boundary value problem, Zakharov equations.

I. INTRODUCTION

Set of coupled nonlinear wave equations describing the interaction between high frequency Langmuir waves and low frequency ion-acoustic waves was first derived by Zakharov [1]. The usual Zakharov system defined in space time \mathbb{R}^{d+1} is expressed as

$$iE_t + \Delta E = nE,$$

 $n_{tt} - \Delta n = \Delta |E|^2$

where $E : \mathbb{R}^{d+1} \to \mathbb{C}^d$ is the slowly varying amplitude of the high-frequency electric field, and $n : \mathbb{R}^{d+1} \to \mathbb{R}$ denotes the fluctuation of the ion-density from its equilibrium.

This system has been the subject of many studies [2], [3], [4], [5], [6], [7], [8], [9], [10]. The 1D Cauchy problem for the Zakharov system is shown to be locally well-posed for low regularity Schrödinger data and wave data under certain assumptions on the parameters [3]. The study [4] deals with the existence and uniqueness of smooth solution for a generalized Zakharov equation and establishes local in time existence and uniqueness in the case of dimensions 2 and 3. S. You studied a generalized Zakharov equation and obtained the existence and uniqueness of the global solutions to the initial value problem [10].

In this paper, we study the following systems of 2dimensional generalized Zakharov type equations

$$iE_t + \Delta E - nE = \alpha |E|^2 E, \qquad (1)$$

$$\partial_t n + \nabla \cdot V = 0, \tag{2}$$

$$V_t + \sum_{i=1}^{2} \frac{\partial \operatorname{grad}\varphi(V)}{\partial x_j} - \beta \Delta V_t + \nabla \left(n + |E|^2 \right) = 0.$$
(3)

with initial boundary data as follows:

$$E|_{t=0} = E_0(x), \quad n|_{t=0} = n_0(x), \quad V|_{t=0} = V_0(x), \quad (4)$$

$$E|_{\partial\Omega} = 0, \quad n|_{\partial\Omega} = 0, \quad V|_{\partial\Omega} = 0$$
 (5)

Manuscript received December 12, 2015; revised October 4, 2016.

Shujun You and Xiaoqi Ning are with the School of Mathematical Sciences, Huaihua University, Huaihua, 418008 China.

E-mail: ysj980@ aliyun.com (S. You), nxq035@163.com (X. Ning).

where $E(x,t) = (E_1(x,t), E_2(x,t), \dots, E_N(x,t))$ is an *N*-dimensional complex valued unknown functional vector, $V(x,t) = (V_1(x,t), V_2(x,t))$ is a 2-dimensional real-valued unknown functional vector, n(x,t) is a real-valued unknown function, $\varphi(s)$ is a real function, and $\beta > 0, x \in \Omega \subset \mathbb{R}^2$, $t \ge 0$.

We study the generalized Zakharov system in dimension 2 with initial boundary data. First, a priori estimates of the problem is made. Next, using the Galerkin method, the existence of the global generalized solution of the problem is shown. In fact, nonlinear partial differential equations have also been studied by others using different approaches, as can be seen in [14]-[36].Now we state the main results of this paper.

Theorem 1. Supposing that

(1)
$$E_0(x) \in H^1(\Omega), V_0(x) \in H^1(\Omega), n_0(x) \in L^2(\Omega),$$

(2) $\varphi(v) \in C^2, \varphi(0) = 0.$
(3) $||E_0(x)||_{L^2}^2 < \frac{\varepsilon}{1+\varepsilon|\alpha|} ||\psi(x)||_{L^2}^2,$
(4) $|\operatorname{grad}\varphi(v)| \leq C (|v|+1).$

where $0 < \varepsilon < 1$, $\psi(x)$ is a solution of the equation

$$\Delta \psi - \psi + \psi^3 = 0.$$

Then there is the global generalized solution of the initial boundary problem (1)-(5).

$$\begin{split} E(x,t) &\in L^{\infty}(\mathbb{R}^{+}; H^{1}) \cap W^{1,\infty}(\mathbb{R}^{+}; H^{-1}), \\ V(x,t) &\in L^{\infty}(\mathbb{R}^{+}; H^{1}) \cap W^{1,\infty}(\mathbb{R}^{+}; H^{1}_{0}), \\ n(x,t) &\in L^{\infty}(\mathbb{R}^{+}; L^{2}) \cap W^{1,\infty}(\mathbb{R}^{+}; L^{2}), \end{split}$$

For the sake of convenience of the following contexts, we set some notations. For $1 \le q \le \infty$, we denote $L^q(\mathbb{R}^d)$ the space of all q times integrable functions in \mathbb{R}^d equipped with norm $\|\cdot\|_{L^{q}(\mathbb{R}^d)}$ or simply $\|\cdot\|_{L^q}$ and $H^{s,p}(\mathbb{R}^d)$ the Sobolev space with norm $\|\cdot\|_{H^{s,p}(\mathbb{R}^d)}$. If p = 2, we write $H^s(\mathbb{R}^d)$ instead of $H^{s,2}(\mathbb{R}^d)$. Let $(f,g) = \int_{\mathbb{R}^n} f(x) \cdot \overline{g(x)} dx$, where $\overline{g(x)}$ denotes the complex conjugate function of g(x). We use C to represent various constants that can depend on the initial data.

This paper is organized as follows. In Section II, we make a priori estimates of the problem (1)-(5). In Section III, we establish the global generalized solution of the problem (1)-(5) by the Galerkin method.

II. A PRIORI ESTIMATES

In this section, we will derive a priori estimates for the solution of the system (1)-(5).

Lemma 1. Suppose that $E_0(x) \in L^2(\Omega)$. Then for the solution of problem (1)-(5), we have

$$||E(x,t)||_{L^{2}(\Omega)}^{2} = ||E_{0}(x)||_{L^{2}(\Omega)}^{2}$$

This work was supported in part by the National Natural Science Foundation of China under Grant 11501232, the Research Foundation of Education Bureau of Hunan Province under Grant 15B185 and 16C1272, and the Scientific Research Found of Huaihua University under Grant HHUY2015-05.

Proof: Taking the inner product of (1) and E, it follows that

$$(iE_t + \Delta E - nE, E) = \left(\alpha |E|^2 E, E\right)$$
(6)

Since

$$\operatorname{Im} (iE_t, E) = \frac{1}{2} \frac{d}{dt} ||E||_{L^2}^2,$$

$$\operatorname{Im} (\Delta E - nE, E) = \operatorname{Im} \left(\alpha |E|^2 E, E \right) = 0,$$

hence from (6), we get

$$\frac{d}{dt} \|E(x,t)\|_{L^2}^2 = 0.$$

We thus get Lemma 1.

Lemma 2. Suppose that

(1)
$$E_0(x) \in H^1(\Omega), V_0(x) \in H^1(\Omega), n_0(x) \in L^2(\Omega),$$

(2) $\varphi(v) \in C^2, \varphi(0) = 0.$

Then we have

$$\mathcal{M}(t) = \mathcal{M}(0).$$

where

$$\mathcal{M}(t) = \|\nabla E\|_{L^2}^2 + \int \left(n |E|^2 + \frac{\alpha}{2} |E|^4\right) dx + \frac{\beta}{2} \|\nabla V\|_{L^2}^2 + \frac{1}{2} \|n\|_{L^2}^2 + \frac{1}{2} \|V\|_{L^2}^2.$$

Proof: Taking the inner products of (1) and E_t , it follows that

$$(iE_t + \Delta E - nE, E_t) = \left(\alpha |E|^2 E, E_t\right). \tag{7}$$

Since

$$\operatorname{Re}\left(iE_{t}, E_{t}\right) = 0,$$

$$\operatorname{Re}\left(\Delta E, E_{t}\right) = -\operatorname{Re}\left(\nabla E, \nabla E_{t}\right) = -\frac{1}{2}\frac{d}{dt}\left\|\nabla E\right\|_{L^{2}}^{2},$$

$$\operatorname{Re}\left(-nE - \alpha |E|^{2} E, E_{t}\right)$$

$$= -\frac{1}{2} \int \left(n\left(|E|^{2}\right)_{t} + \alpha |E|^{2}\left(|E|^{2}\right)_{t}\right) dx$$

$$= -\frac{1}{2} \frac{d}{dt} \int \left(n|E|^{2} + \frac{\alpha}{2}|E|^{4}\right) dx + \frac{1}{2} \int n_{t}|E|^{2} dx.$$
(7)

thus from (7) it follows that

$$\frac{d}{dt}\left(\|\nabla E\|_{L^{2}}^{2} + \int \left(n|E|^{2} + \frac{\alpha}{2}|E|^{4}\right)dx\right) = \int n_{t}|E|^{2}dx.$$
 (8)

From (2)-(3), we obtain

$$\int n_{t}|E|^{2}dx = -\int (\nabla \cdot V)|E|^{2}dx = \int V \cdot \nabla |E|^{2}dx$$

$$= \int V \cdot \left(-V_{t} - \sum_{j=1}^{2} \frac{\partial \operatorname{grad}\varphi(V)}{\partial x_{j}} + \beta \Delta V_{t} - \nabla n\right) dx$$

$$= -\beta \int \nabla V \cdot \nabla V_{t}dx + \int (\nabla \cdot V) n dx - \frac{1}{2} \frac{d}{dt} ||V||_{L^{2}}^{2}$$

$$+ \sum_{j=1}^{2} \int \operatorname{grad}\varphi(V) \cdot \frac{\partial V}{\partial x_{j}} dx$$

$$= -\frac{\beta}{2} \frac{d}{dt} ||\nabla V||_{L^{2}}^{2} - \int n_{t} n dx - \frac{1}{2} \frac{d}{dt} ||V||_{L^{2}}^{2} + \sum_{j=1}^{2} \int \frac{\partial \varphi(V)}{\partial x_{j}} dx$$

$$= -\frac{1}{2} \frac{d}{dt} \left(\beta ||\nabla V||_{L^{2}}^{2} + ||n||_{L^{2}}^{2} + ||V||_{L^{2}}^{2} \right). \tag{9}$$

Combining (8) with (9) we obtain

$$\frac{d}{dt} \left(\|\nabla E\|_{L^2}^2 + \int \left(n|E|^2 + \frac{\alpha}{2}|E|^4 \right) dx \right) \\ + \frac{1}{2} \frac{d}{dt} \left(\beta \|\nabla V\|_{L^2}^2 + \|n\|_{L^2}^2 + \|V\|_{L^2}^2 \right) = 0.$$

Letting

$$\begin{aligned} \mathcal{M}(t) &= \left\| \nabla E \right\|_{L^2}^2 + \int \left(n |E|^2 + \frac{\alpha}{2} |E|^4 \right) dx \\ &+ \frac{1}{2} \left(\beta \| \nabla V \|_{L^2}^2 + \| n \|_{L^2}^2 + \| V \|_{L^2}^2 \right). \end{aligned}$$

It follows that

Lemma 3 (Gagliardo-Nirenberg inequality [11]). Assume that $u \in L^q(\mathbb{R}^n)$, $D^m u \in L^r(\mathbb{R}^n)$, $1 \le q, r \le \infty, 0 \le j \le m$, we have the estimations

 $\mathcal{M}(t) = \mathcal{M}(t).$

$$|D^{j}u||_{L^{p}(\mathbb{R}^{n})} \leq C ||D^{m}u||_{L^{r}(\mathbb{R}^{n})}^{\alpha}||u||_{L^{q}(\mathbb{R}^{n})}^{1-\alpha}$$

where C is a positive constant, $0 \le \frac{j}{m} \le \alpha \le 1$,

$$\frac{1}{p} = \frac{j}{n} + \alpha \left(\frac{1}{r} - \frac{m}{n} \right) + (1 - \alpha) \frac{1}{q}.$$

Lemma 4 (Sobolev's best constant estimates [12]). Suppose that $f(x) \in H^1(\mathbb{R}^N)$. Then we have

$$\begin{split} \|f\|_{L^{2p+2}(\mathbb{R}^N)}^{2p+2} &\leq C_{p,N}^{2p+2} \|\nabla f\|_{L^2(\mathbb{R}^N)}^{pN} \|f\|_{L^2(\mathbb{R}^N)}^{2+p(2-N)}, \\ 0 &$$

where the constant

$$C_{p,N} = \left(\frac{p+1}{\|\psi\|_{L^2(\mathbb{R}^N)}^{2p}}\right)^{\frac{1}{2p+1}}$$

and $\psi(x)$ is a ground state solution for the equation

$$\frac{pN}{2}\Delta\psi - \left(1 + \frac{p}{2}(2 - N)\right)\psi + \psi^{2p+1} = 0.$$
 (10)

Obviously, the solution of equation (10) exists, and $\psi(x) \neq 0$ *.*

Lemma 5. Suppose that the conditions of Lemma 2 are satisfied and

$$\|E_0(x)\|_{L^2}^2 < \frac{\varepsilon}{1+\varepsilon |\alpha|} \, \|\psi(x)\|_{L^2}^2 \, ,$$

where $0 < \varepsilon < 1$, ψ is a solution of the equation

$$\Delta\psi - \psi + \psi^3 = 0.$$

Then we have

$$|\nabla E||_2^2 + ||\nabla V||_2^2 + ||n||_2^2 + ||V||_2^2 \le C.$$

Proof: By Hölder inequality, Young inequality, there holds

$$\begin{split} \int_{\Omega} n |E|^2 dx \bigg| &\leq \|n\|_{L^2} \|E\|_{L^4}^2 \\ &\leq \frac{\varepsilon}{2} \|n\|_{L^2}^2 + \frac{1}{2\varepsilon} \|E\|_{L^4}^4. \end{split}$$
(11)

Using Gagliardo-Nirenberg inequality and Lemma 4, we express as

$$\|E\|_{L^4}^4 \le \frac{2}{\|\psi\|_{L^2}^2} \|\nabla E\|_{L^2}^2 \|E\|_{L^2}^2.$$
(12)

(Advance online publication: 24 May 2017)

Note that from Lemma 2 and Equations (11), (12), one has

$$\left(1 - \frac{1 + \varepsilon |\alpha|}{\varepsilon ||\psi||_{L^2}^2} ||E_0||_{L^2}^2 \right) ||\nabla E||_{L^2}^2 + \frac{\beta}{2} ||\nabla V||_{L^2}^2 + \frac{1 - \varepsilon}{2} ||n||_{L^2}^2 + \frac{1}{2} ||V||_{L^2}^2 \le |\mathcal{M}(0)|.$$

Note that

$$\left\|E_0(x)\right\|_{L^2}^2 < \frac{\varepsilon}{1+\varepsilon|\alpha|} \left\|\psi(x)\right\|_L^2$$

and $0 < \varepsilon < 1$, we thus get Lemma 5.

Lemma 6. Suppose that the conditions of Lemma 5 are satisfied and $|\operatorname{grad}\varphi(v)| \leq C(|v|+1)$. Then we have

$$||E_t||_{H^{-1}} + ||V_t||_{H^1_0} + ||n_t||_{L^2} \le C.$$

Proof: Taking the inner product of Eq. (1) and Φ , Eq. (2) and φ , it follows that

$$(iE_t + \Delta E - nE, \Phi) = \left(\alpha |E|^2 E, \Phi\right), \tag{13}$$

$$(\partial_t n + \nabla \cdot V, \varphi) = 0, \tag{14}$$

where $\varphi, \varphi_j \in H_0^2$ $(j = 1, \dots, N), \Phi = (\varphi_1, \dots, \varphi_N).$ By Hölder inequality, it follows from Eq. (13) that

$$\begin{aligned} |(E_{t}, \Phi)| &\leq |(\Delta E, \Phi)| + |(nE, \Phi)| + \left| \left(\alpha |E|^{2}E, \Phi \right) \right| \\ &= |(\nabla E, \nabla \Phi)| + |(nE, \Phi)| + \left| \left(\alpha |E|^{2}E, \Phi \right) \right| \\ &\leq ||\nabla E||_{L^{2}} \, ||\nabla \Phi||_{L^{2}} + ||n||_{L^{2}} \, ||E||_{L^{4}} \, ||\Phi||_{L^{4}} \\ &+ |\alpha| \, ||E||_{L^{6}}^{3} \, ||\Phi||_{L^{2}} \, . \end{aligned}$$
(15)

By Gagliardo-Nirenberg inequality, we know that

$$\begin{split} \|E\|_{L^4} &\leq C \|\nabla E\|_{L^2}^{\frac{1}{2}} \|E\|_{L^2}^{\frac{1}{2}} \leq C, \\ \|E\|_{l^6} &\leq C \|\nabla E\|_{L^2}^{\frac{2}{3}} \|E\|_{L^2}^{\frac{1}{3}} \leq C, \\ \|\Phi\|_{L^4} &\leq C \|\nabla \Phi\|_{L^2}^{\frac{1}{2}} \|\Phi\|_{L^2}^{\frac{1}{2}} \leq C \left(\|\nabla \Phi\|_{L^2} + \|\Phi\|_{L^2} \right). \end{split}$$

Hence from Eq. (15) we get

$$|(E_t, \Phi)| \le C ||\Phi||_{H^1_0}.$$
 (16)

Using Hölder inequality, from Eq. (14), there is

$$|(n_{t},\varphi)| = |(\nabla \cdot V,\varphi)| \le ||\nabla V||_{L^{2}} \, ||\varphi||_{L^{2}} \le C ||\varphi||_{L^{2}}.$$
 (17)

Taking the inner product of Eq. (3) and V_t , it follows that

$$\left(V_t + \sum_{j=1}^2 \frac{\partial \operatorname{grad}\varphi(V)}{\partial x_j} - \beta \Delta V_t + \nabla \left(n + |E|^2\right), V_t\right) = 0. \quad (18)$$

Since

$$(V_t - \beta \Delta V_t, V_t) = ||V_t||_{L^2}^2 + \beta ||\nabla V_t||_{L^2}^2,$$

$$\left| \left(\sum_{j=1}^{2} \frac{\partial \operatorname{grad} \varphi(V)}{\partial x_{j}}, V_{t} \right) \right| = \left| \left(\sum_{j=1}^{2} \operatorname{grad} \varphi(V), \frac{\partial V_{t}}{\partial x_{j}} \right) \right|$$

$$\leq C \left(||V||_{L^{2}} + 1 \right) ||\nabla V_{t}||_{L^{2}}$$

$$\leq C + \frac{\beta}{4} ||\nabla V_{t}||_{L^{2}}^{2},$$

 $|(\nabla n, V_t)| = |(n, \nabla \cdot V_t)| \le ||n||_{L^2} ||\nabla V_t||_{L^2} \le C + \frac{\beta}{4} ||\nabla V_t||_{L^2}^2,$

$$\begin{split} \left| \left(\nabla |E|^2, V_t \right) \right| &= \left| \left(|E|^2, \nabla \cdot V_t \right) \right| \le ||E||_{L^4}^2 ||\nabla V_t||_{L^2} \\ &\le C + \frac{\beta}{4} ||\nabla V_t||_{L^2}^2. \end{split}$$

from Eq. (18) we get

$$\|V_t\|_{L^2}^2 + \frac{\beta}{4} \|\nabla V_t\|_{L^2}^2 \le C.$$
(19)

Hence from (16), (17) and (19), we obtain

$$||E_t||_{H^{-1}} + ||V_t||_{H^1_0} + ||n_t||_{L^2} \le C.$$

In this section, we formulate the proof of Theorem 1. First we give the definition of generalized solution for problems (1)-(5).

Definition 1. The functions

$$\begin{split} E_m(x,t) &\in L^{\infty}(\mathbb{R}^+; H^1) \cap W^{1,\infty}(\mathbb{R}^+; H^{-1}), \quad m = 1, 2, \cdots, N, \\ V_{\lambda}(x,t) &\in L^{\infty}(\mathbb{R}^+; H^1) \cap W^{1,\infty}(\mathbb{R}^+; H^1_0), \quad \lambda = 1, 2, \\ n(x,t) &\in L^{\infty}(\mathbb{R}^+; L^2) \cap W^{1,\infty}(\mathbb{R}^+; L^2) \end{split}$$

are called generalized solution of problems (1)-(5), if for any $\xi \in H_0^1$ they satisfy the integral equality

$$(iE_{mt},\xi) - (\nabla E_m, \nabla \xi) - (nE_m,\xi) = (\alpha |E|^2 E_m,\xi),$$

$$m = 1, 2, \cdots, N,$$

$$(\partial_t n, \xi) - (V, \nabla \xi) = 0,$$

$$(V_{\lambda t}, \xi) - \left(\sum_{j=1}^2 \frac{\partial \varphi(V)}{\partial V_\lambda}, \frac{\partial \xi}{\partial x_j}\right) + (\beta \nabla V_{\lambda t}, \nabla \xi)$$

$$- \left(n + |E|^2, \frac{\partial \xi}{\partial x_\lambda}\right) = 0, \ \lambda = 1, 2.$$

with initial boundary data as follows:

$$\begin{split} E|_{t=0} &= E_0(x), \quad n|_{t=0} = n_0(x), \quad V|_{t=0} = V_0(x), \\ E|_{\partial\Omega} &= 0, \quad n|_{\partial\Omega} = 0, \quad V|_{\partial\Omega} = 0 \end{split}$$

Next, we give two lemmas recalled in [13].

Lemma 7. Let B_0, B, B_1 be three reflexive Banach spaces and assume that the embedding $B_0 \rightarrow B$ is compact. Let

$$W = \left\{ V \in L^{p_0}((0,T); B_0), \ \frac{\partial V}{\partial t} \in L^{p_1}((0,T); B_1) \right\},$$
$$T < \infty, \ 1 < p_0, p_1 < \infty.$$

W is a Banach space with norm

$$||V||_{W} = ||V||_{L^{p_{0}}((0,T);B_{0})} + ||V_{t}||_{L^{p_{1}}((0,T);B_{1})}.$$

Then the embedding $W \to L^{p_0}((0,T); B)$ is compact.

Lemma 8. Let Ω be an open set of \mathbb{R}^n and let $g, g_{\varepsilon} \in$ $L^p(\mathbb{R}^n), 1 , such that$

$$g_{\varepsilon} \to g$$
 a.e. in Ω and $||g_{\varepsilon}||_{L^{p}(\Omega)} \leq C$.

Then $g_{\varepsilon} \to g$ weakly in $L^{p}(\Omega)$.

Now, one can estimate Theorem 1.

Proof: By using the Galerkin method, choose the basic periodic functions $\{\omega_i(x)\}$ as follows:

$$-\Delta\omega_j(x) = \lambda_j\omega_j(x), \ \omega_j(x) \in H^1_0(\Omega), \ j = 1, 2, \cdots, l.$$

(Advance online publication: 24 May 2017)

The approximate solution of problem (1)-(5) can be written It then follows from Eq. (27) and (29) that as

$$\begin{split} E^{l}(x,t) &= \sum_{j=1}^{l} \alpha_{j}^{l}(t)\omega_{j}(x), \quad V^{l}(x,t) = \sum_{j=1}^{l} \beta_{j}^{l}(t)\omega_{j}(x), \\ n^{l}(x,t) &= \sum_{j=1}^{l} \gamma_{j}^{l}(t)\omega_{j}(x), \end{split}$$

where

$$E^{l} = \left(E_{1}^{l}, \cdots, E_{N}^{l}\right), \quad \alpha_{j}^{l}(t) = \left(\alpha_{j1}^{l}(t), \cdots, \alpha_{jN}^{l}(t)\right),$$
$$V^{l} = \left(V_{1}^{l}, V_{2}^{l}\right), \quad \beta_{j}^{l}(t) = \left(\beta_{j1}^{l}(t), \beta_{j2}^{l}(t)\right).$$

According to Galerkin's method, these undetermined coefficients $\alpha_i^l(t)$, $\beta_i^l(t)$ and $\gamma_i^l(t)$ must satisfy the following initial value problem of the system of ordinary differential equations:

$$\begin{aligned} \left(iE_{ml}^{l},\omega_{\kappa}\right) &- \left(\nabla E_{m}^{l},\nabla \omega_{\kappa}\right) - \left(n^{l}E_{m}^{l},\omega_{\kappa}\right) \\ &= \left(\alpha\left|E^{l}\right|^{2}E_{m}^{l},\omega_{\kappa}\right), \quad m = 1, 2, \cdots, N, \end{aligned}$$
 (20)

$$\left(\partial_t n^l, \omega_\kappa\right) - \left(V^l, \nabla \omega_\kappa\right) = 0, \quad \kappa = 1, 2, \cdots, l,$$
 (21)

$$\begin{pmatrix} V_{\lambda l}^{l}, \omega_{\kappa} \end{pmatrix} - \left(\sum_{j=1}^{2} \frac{\partial \varphi(V^{l})}{\partial V_{\lambda}^{l}}, \frac{\partial \omega_{\kappa}}{\partial x_{j}} \right) + \left(\beta \nabla V_{\lambda l}^{l}, \nabla \omega_{\kappa} \right) - \left(n^{l} + \left| E^{l} \right|^{2}, \nabla \omega_{\kappa} \right) = 0, \quad \lambda = 1, 2.$$
 (22)

with initial boundary data:

$$E^{l}|_{t=0} = E_{0}(x), \quad n^{l}|_{t=0} = n_{0}(x), \quad V^{l}|_{t=0} = V_{0}(x),$$
 (23)

$$\mathcal{E}^l|_{\partial\Omega} = 0, \quad n^l|_{\partial\Omega} = 0, \quad V^l|_{\partial\Omega} = 0$$
 (24)

Supposing

1

$$E_0^l(x) \xrightarrow{H^1} E_0(x), \quad V_0^l(x) \xrightarrow{H^1} V_0(x),$$
$$n_0^l(x) \xrightarrow{L^2} n_0(x), \quad l \to \infty.$$

Similar to the proof of Lemma 1-6, for the solution $E^{l}(x, t)$, $V^{l}(x, t)$ and $n^{l}(x, t)$ of problem (20)-(24), we can establish the following estimations:

$$\left\|E^{l}\right\|_{H^{1}}+\left\|V^{l}\right\|_{H^{1}}+\left\|n^{l}\right\|_{L^{2}}\leq C$$
(25)

$$\left\|E_{t}^{l}\right\|_{H^{-1}} + \left\|V_{t}^{l}\right\|_{H^{1}_{0}} + \left\|n_{t}^{l}\right\|_{L^{2}} \le C$$
(26)

where the constant C is independent of l and D. By compact argument, some subsequence of (E^l, V^l, n^l) , also labeled as *l*, has a weak limit (E, V, n). More precisely,

$$E^{l}(x,t) \to E(x,t)$$
 in $L^{\infty}(\mathbb{R}^{+};H^{1})$ weakly star, (27)

$$V^{l}(x,t) \to V(x,t)$$
 in $L^{\infty}(\mathbb{R}^{+}; H^{1})$ weakly star, (28)

$$n^{l}(x,t) \to n(x,t)$$
 in $L^{\infty}(\mathbb{R}^{+};L^{2})$ weakly star. (29)

Eq. (26) implies that

$$E_t^l \to E_t \text{ in } L^{\infty}(\mathbb{R}^+; H^{-1}) \text{ weakly star,}$$
(30)
$$V_t^l \to V_t \text{ in } L^{\infty}(\mathbb{R}^+; H_0^1) \text{ weakly star,}$$

 $n_t^l \to n_t$ in $L^{\infty}(\mathbb{R}^+; L^2)$ weakly star.

Moreover, it should be noted that the following maps are continuous.

$$\begin{split} H^1 &\to L^4, \quad u \mapsto u, \\ H^1 &\times L^2 \to L^2, \quad (u,v) \mapsto uv. \end{split}$$

$$|E^l|^2 \to w \text{ in } L^{\infty}(\mathbb{R}^+; L^2) \text{ weakly star,}$$
 (31)

$$n^{l}E^{l} \to z \text{ in } L^{\infty}(\mathbb{R}^{+}; L^{2}) \text{ weakly star.}$$
 (32)

First, we prove $w = |E|^2$. Let Π be any bounded subdomain of Ω . We notice that

the embedding $H^1(\Pi) \to L^4(\Pi)$ is compact.

and for any Banach space X,

the embedding $L^{\infty}(\mathbb{R}^+; X) \to L^2(0, T; X)$ is continuous.

Hence, according to Eq. (27), (31) and Lemma 7, applied to $B_0 = H^1(\Pi), B = L^4(\Pi), B_1 = H^{-1}(\Pi)$, and says that some subsequence of $E^l|_{\Pi}$ (also referred to as *l*) converges strongly to $E|_{\Pi}$ in $L^2(0,T;L^4(\Pi))$. So we can assume that

$$E^l \to E$$
 strongly in $L^2(0, T; L^4_{loc}(\Pi))$,

and thus

$$E^l \to E$$
 a.e. in $[0, T] \times \Pi$.

Then, using Lemma 8 and Eq. (31) imply that $w = |E|^2$ Second, we prove z = nE. Let χ be a test function in $L^2(0,T;H^1)$, supp $\chi \subset \Omega \subset \mathbb{R}^2$.

$$\int_0^T \int_\Omega (n^l E^l - nE) \chi dx dt$$

=
$$\int_0^T \int_\Omega n^l (E^l - E) \chi dx dt + \int_0^T \int_\Omega (n^l - n) E \chi dx dt.$$

On one hand

$$\begin{split} & \left| \int_{0}^{T} \int_{\Omega} n^{l} \left(E^{l} - E \right) \chi dx dt \right| \\ & \leq \left\| n^{l} \right\|_{L^{\infty}(0,T;L^{2}(\Omega))} \left\| E^{l} - E \right\|_{L^{2}(0,T;L^{4}(\Omega))} \left\| \chi \right\|_{L^{2}(0,T;L^{4}(\Omega))}, \end{split}$$

Since Ω is bounded, we deduce from Eq. (27) and (29) that

$$\left|\int_0^T \int_{\Omega} n^l \left(E^l - E\right) \chi dx dt\right| \to 0 \quad (l \to +\infty).$$

On the other hand, it should be noted that $E\chi \in L^1(0, T; L^2)$. In fact,

$$\|E\chi\|_{L^1(0,T;L^2)} \le \|E\|_{L^2(0,T;L^4)} \, \|\chi\|_{L^2(0,T;L^4)} < \infty.$$

Therefore we deduce from Eq. (29) that

$$\int_0^T \int_\Omega \left(n^l - n \right) E \chi dx dt \to 0 \quad (l \to +\infty).$$

Thus $n^l E^l \rightarrow nE$ in $L^2(o, T; H^{-1})$. So z = nE.

Hence taking $l \rightarrow \infty$ from Eq. (20)-(24), by using the density of ω_i in $H_0^1(\Omega)$ we get the existence of local generalized solution for the periodic initial boundary value problem (1)-(5); letting $D \to \infty$, the existence of local solution for the initial boundary value problem (1)-(5) can be obtained. By the continuation extension principle and a priori estimates, we can get the existence of a global generalized solution for problem (1)-(5).

We thus complete the proof of Theorem 1.

(Advance online publication: 24 May 2017)

References

- V.E. Zakharov, "Collapse of Langmuir waves," Sov. Phys. JETP, vol. 35, pp. 908-914, 1972.
- [2] J. Holmer, "Local ill-posedness of the 1D Zakharov system," *Electron. J. Differential Equations*, vol. 24, pp. 1-24, 2007.
- [3] H. Pecher, "An improved local well-posedness result for the onedimensional Zakharov system," J. Math. Anal. Appl., vol. 342, pp. 1440-1454, 2008.
- [4] B. Guo, J. Zhang and X. Pu, "On the existence and uniqueness of smooth solution for a generalized Zakharov equation," J. Math. Anal. Appl., vol. 365, pp. 238-253, 2010.
- [5] F. Linares and C. Matheus, "Well Posedness for the 1D Zakharov-Rubenchik system," Advances in Differential Equations, vol. 14, pp. 261-288, 2009.
- [6] F. Linares and J.C. Saut, "The Cauchy problem for the 3D Zakharov-Kuznetsov equation," *Discrete and Continuous Dynamical Systems*, vol. 24, pp. 547-565, 2009.
- [7] Dem zo and Jahrestag der DDR gewidmet, "Über N\u00e4herungsverfahren zur L\u00f6sung der nichtlinearen Schr\u00f6dinger-Gleichung mit selbstkonsistentem Potential," *Math. Nachr*, vol. 89, no. 1, pp. 71-85, 1979.
- [8] S. You, B. Guo and X. Ning, "Equations of Langmuir Turbulence and Zakharov Equations£Smoothness and Approximation," *Appl. Math. Mech.*, vol. 33, no. 8, pp. 1079-1092, 2012.
- [9] N. Masmoudi and K. Nakanishi, "Energy convergence for singular limits of Zakharov type systems," *Invent. Math*, vol. 172, pp. 535-583, 2008.
- [10] S. You, "The posedness of the periodic initial value problem for generalized Zakharov equations," *Nonlinear Analysis: TMA*, vol. 71, no. 7-8, pp. 3571-3584, 2009.
- [11] A. Friedman, *Partial Differential Equations*, Holt, Reihart and Winston, Inc., New York, Montreal Que, London, 1969.
- [12] M.I. Weinstein, "Nonlinear Schrödinger equations and sharp interpolation estimates," *Commun Math Phys*, vol. 87, no. 4, pp. 567-576, 1983.
- [13] J.L. Lions, *Quelques methods de resolution des problemes aux limites non lineaires*, Dunod Gauthier Villard, Paris, 1969.
- [14] K. Mohammed, "New Exact traveling wave solutions of the (3+1) dimensional Kadomtsev-Petviashvili (KP) equation," *IAENG International Journal of Applied Mathematics*, vol. 37, no. 1, pp. 17-19, 2007.
- [15] N. Shang and B. Zheng, "Exact solutions for three fractional partial differential equations by the (G/G) method," *IAENG International Journal of Applied Mathematics*, vol. 43, no. 3, pp. 114-119, 2013.
- [16] Y. Huang, "Explicit multi-soliton solutions for the KdV equation by Darboux transformation," *IAENG International Journal of Applied Mathematics*, vol. 43, no. 3, pp. 135-137, 2013.
- [17] A. El-Ajou, Z. Odibat and S. Momani et al, "Construction of analytical solutions to fractional differential equations using homotopy analysis method," *IAENG International Journal of Applied Mathematics*, vol. 40, no. 2, pp. 43-51, 2010.
- [18] B.S. Desale and G.K. Srinivasan, "Singular analysis of the system of ODE reductions of the stratified Boussinesq equations," *IAENG International Journal of Applied Mathematics*, vol. 38, no. 4, pp. 184-191, 2008.
- [19] J.M. Yoon, S. Xie and V. Hrynkiv, "A series solution to a partial integro-differential equation arising in viscoelasticity," *IAENG International Journal of Applied Mathematics*, vol. 43, no. 4, pp. 172-175, 2013.
- [20] C. Yang and J. Hou, "Numerical method for solving Volterra integral equations with a convolution kernel," *IAENG International Journal of Applied Mathematics*, vol. 43, no. 4, pp. 185-189, 2013.
- [21] D.E. Panayotounakos, T.I. Zarmpoutis and P. Sotiropoulos, "The general solutions of the normal Abels type nonlinear ODE of the second kind," *IAENG International Journal of Applied Mathematics*, vol. 43, no. 3, pp. 94-98, 2013.
- [22] F. Samat, F. Ismail and M. Suleiman, "Phase fitted and amplification fitted hybrid methods for solving second-order ordinary differential equations," *IAENG International Journal of Applied Mathematics*, vol. 43, no. 3, pp. 99-105, 2013.
- [23] N.A. Larkin and E. Tronco, "Regular solutions of the 2D Zakharov-Kuznetsov equation on a half-strip," *J. Differential Equations*, vol. 254, no. 1, pp. 81-101, 2013.
- [24] I. Aslan, "Generalized solitary and periodic wave solutions to a (2 +1)-dimensional ZakharovCKuznetsov equation," *Applied Mathematics and Computation*, vol. 217, no. 4, pp. 1421-1429, 2010.
- [25] I. Aslan, "The frist integral method for constructing exact and explicit solutions to nonlinear evolution equations," *Math. Meth. Appl. Sci.*, vol. 35, no. 6, pp. 716-722, 2012.
- [26] A.R. Seadawy, "Stability analysis for ZakharovCKuznetsov equation of weakly nonlinear ion-acoustic waves in a plasma," *Computers and Mathematics with Applications*, vol. 67, no. 1, pp. 172-180, 2014.

- [27] N.A. Larkin and E. Tronco, "Nonlinear quarter-plane problem for the Korteweg-de Vries equation," J. Differential Equations, vol. 2011, no. 113, pp. 1-22, 2011.
- [28] F. Linares and A. Pastor, "Well-posedness for the 2D modified Zakharov-Kuznetsov equation," J. Funct. Anal., vol. 260, no. 4, pp. 1060-1085, 2011.
- [29] F. Linares, A. Pastor and J.C. Saut, "Well-posedness for the ZK equation in a cylinder and on the background of a KdV soliton," *Comm. Partial Differential Equations*, vol. 35, no. 9, pp. 1674-1689, 2010.
- [30] F. Linares and J.C. Saut, "The Cauchy problem for the 3D Zakharov-Kuznetsov equation," *Discrete Contin. Dyn. Syst.*, vol. 24, pp. 547-565, 2009.
- [31] F. Linares and A.F. Pazoto, "Asymptotic behavior of the Korteweg-de Vries equation posed in a quarter plane," *J. Differential Equations*, vol. 246, no. 4, pp. 1342-1353, 2009.
 [32] I. Rivas, M. Usman and B. Zhang, "Global well-posedness and
- [32] I. Rivas, M. Usman and B. Zhang, "Global well-posedness and asymptotic behavior of a class of initial-boundary value problem for the Korteweg-de Vries equation on a finite domain," *Math. Control Relat. Fields*, vol. 1, no. 1, pp. 61-81, 2011.
- [33] L. Rosier, "Exact boundary controllability for the Korteweg-de Vries equation on a bounded domain," *ESAIM Control Optim. Calc. Var.*, vol. 2, pp. 33-55, 1997.
- [34] L. Rosier, "A survey of controllability and stabilization results for partial differential equations," *RS-JESA*, vol. 41, no. 3-4, pp. 365-411, 2007.
- [35] L. Rosier and B. Zhang, "Control and stabilization of the KdV equation: recent progress," J. Syst. Sci. Complex., vol. 22, no. 4, pp. 647-682, 2009.
- [36] J.C. Saut and R. Temam, "An initial boundary-value problem for the Zakharov-Kuznetsov equation," *Adv. Differential Equations*, vol. 15, no. 11-12, pp. 1001-1031, 2010.