# A Prior Error Estimate for Linear Finite Element Approximation to Interface Optimal Control Problems 

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#### Abstract

This paper considers a linear finite element method for the constrained optimal control problems (OCPs) governed by elliptic interface equations. The state and adjoint state are approximated by the conforming $P_{1}$ elements, while the control is approximated with the orthogonal projection of the adjoint state. Optimal order error estimates are proved in both $L^{2}$ norm and broken energy norm. Lastly, some numerical results are presented to confirm the theoretical analysis.


Index Terms-finite element method, interface OCPs, optimal order error estimates.

## I. Introduction

0Ptimal control problems (OCPs) governed partial differential equations are playing an increasingly crucial role in a lot of engineering applications, such as chemical processes, fluid dynamics, medicine, economics, and so on [3], [24]. Much attention has been paid to the numerical solution of these problems since their analytical solutions do not always exist. In the recent decades, finite element methods (FEM) have been developed to be one of the most popular and efficient methods not only for partial differential equations [26], but also for many scientific computing fields, i.e., the magnetic resonance elastography [18], mechanism analysis [19], predicting the blasting effect [27], etc. Recently, FEMs have been intensively investigated for OCPs governed by partial differential equations. A priori error estimate was firstly proposed in [6] for the OCPs and obtained the error estimates in $L^{2}$-norm. [14] derived the error estimates of FEM for an elliptic OCPs with a small parameter. Mixed FEM for OCPs governed by elliptic equations and Stokes equations was presented in [4] and [15]. On the other hand, some a posteriori error estimates of conforming FEMs for the OCPs were reported in [12], [13], [21] and the references cited therein. In addition, some

[^0]discussions on nonconforming FEMs for OCPs can be found from [7], [8], [10], [11].

We consider the following interface OCPs: find $(y, u) \in$ $Y \times U$, such that

$$
\begin{equation*}
\min _{u \in U_{a d} \subset U} J(y, u)=\frac{1}{2}\left\|y-y_{d}\right\|_{0, \Omega}^{2}+\frac{\alpha}{2}\|u\|_{0, \Omega}^{2} \tag{1}
\end{equation*}
$$

subject to

$$
\left\{\begin{array}{l}
-\nabla \cdot(\beta \nabla y)=u, \text { in } \Omega  \tag{2}\\
y=0, \text { on } \partial \Omega \\
{[y]_{\Gamma}=0,\left[\beta \frac{\partial y}{\partial n}\right]_{\Gamma}=0,}
\end{array}\right.
$$

where $\alpha$ is a positive constant parameter, $\Omega$ is a convex polygon in $R^{2}$. Let $\Omega^{-} \subset \Omega$ be an open domain with a $C^{2}$ curve boundary $\Gamma \subset \Omega$, and let $\Omega^{+}=\Omega \backslash \Omega^{-}$ (see Fig.1). Throughout this paper, we use the standard


Fig. 1. $\Omega=\Omega^{-} \cup \Omega^{+}$

Sobolev spaces and norms (see [2]), and further denote $Y=H_{0}^{1}(\Omega) \cap H^{2}\left(\Omega^{-}\right) \cap H^{2}\left(\Omega^{+}\right)$, and $U=H^{1}(\Omega)$. The target state $y_{d} \in C^{0}(\bar{\Omega})$ is a given function. The admissible control set $U_{a d}$ is defined as

$$
\begin{equation*}
U_{a d}=\{v \in U: a(x) \leq v \leq b(x), \text { a.e. in } \Omega\} \tag{3}
\end{equation*}
$$

in which $a(x), b(x) \in L^{\infty}(\Omega)$, and $a(x)<b(x)$.
In (2), we denote by $[v]_{\Gamma}$ the jump of $v$ across the interface $\Gamma$ and $n$ the unit outward normal to $\Gamma$, respectively. The coefficient $\beta$ is a positive piecewise constant function defined by

$$
\begin{equation*}
\beta(x)=\beta^{s}, \quad x \in \Omega^{s}, \tag{4}
\end{equation*}
$$

where $s=-$ or + .
The interface OCP (1)-(2) has remarkable application backgrounds, such as the optimization or optimal control of a process in a domain which is composed of several materials separated by interfaces. Coefficients in partial differential equations may have a jump across the interface among different materials. There are mainly two approaches for numerically solving interface OCPs by using FEM. The first one is to utilize conventional FEMs or its variations
defined on a body-fitted mesh for the domain that contains a interface [1]. Another approach that has drawn more attention recently is the so-called immersed FEM [16], [17], [25]. This method constructs a finite element space that allows piecewise continuous basis functions on each element in order to approximate the interface jump conditions.
In this paper, we present a $P_{1}$-conforming triangular bodyfitted FEM approximation to the elliptic interface OCP (1)(2), which could also be extended to parabolic and hyperbolic OCPs. This method was studied in [5] for solving interface problems and obtained the suboptimal order error estimates in $H^{1}$ and $L^{2}$ norms when the interface is of $C^{2}$ smooth. The authors of this paper also pointed out that the error estimate in $H^{1}$ norm can be optimal if the exact solution belongs to $W^{1, \infty}$ near the interface (cf. Remark 2.4 in [5]). Later on, [23] provided the detailed proof of the above statement, and [9] extended this method to $P_{1}$-nonconforming element.
The remainder of this paper is organized as follows: in Section II, we present the discrete formulations and some useful lemmas. Then, in Section III, we derive the optimal order error estimates for both the state variable and the control variable. In the last section, some numerical results are given to verify the validity of the proposed method.

## II. The discrete formulation and some lemmas

We know from [20] that (1)-(2) has a unique solution $(y, u)$ if and only if there is an adjoint state $p \in Y$, such that $(y, p, u)$ satisfies the following optimality conditions:

$$
\left\{\begin{array}{l}
a(y, v)=(u, v), \quad \forall v \in Y  \tag{5}\\
a(p, v)=\left(y-y_{d}, v\right), \quad \forall v \in Y \\
(\alpha u+p, v-u) \geq 0, \quad \forall v \in U_{a d}
\end{array}\right.
$$

where $a(y, v)=\int_{\Omega} \beta \nabla y \nabla v d x ;(u, v)=\int_{\Omega} u v d x ; p \in Y$ is the adjoint state variables. Specifically, the second equation of (5) is the weak form of

$$
\left\{\begin{array}{l}
-\nabla \cdot(\beta \nabla p)=y-y_{d}, \quad \text { in } \Omega,  \tag{6}\\
p=0, \quad \text { on } \partial \Omega, \\
{[p]_{\Gamma}=0, \quad\left[\beta \frac{\partial p}{\partial n}\right]_{\Gamma}=0,}
\end{array}\right.
$$

where $p$ also satisfies the jump condition as same as $y$ for it in (1).

In addition, with the admissible control set (3), we can get the explicit representation of the optimal control $u$ through the adjoint state $p$,

$$
\begin{align*}
u(x) & =P_{U_{a d}}\left\{-\frac{1}{\alpha} p(x)\right\} \\
& =\min \left\{b(x), \max \left(a(x),-\frac{1}{\alpha} p(x)\right)\right\} \tag{7}
\end{align*}
$$

in which $P_{U_{a d}}$ denotes the orthogonal projection operator onto $U_{a d}$.
Next, we introduce a quasi-uniform triangulation $T_{h}=$ $\{K\}$ of the domain $\Omega$ as in [5], [9]. We denote the diameter of $K$ by $h_{K}$, and let $h=\max _{K \in T_{h}} h_{K}$.

To decompose the interface $\Gamma$, we first approximate the domain $\Omega^{-}$by a region $\Omega_{h}^{-}$with a polygonal boundary $\Gamma_{h}$ whose vertices all lie on the interface $\Gamma$. Let $\Omega_{h}^{+}=\Omega-\Omega_{h}^{-}$. Then, we require each $K \in T_{h}$ to satisfy the following two conditions (see Fig. 2):
(i) $K$ is either in $\Omega_{h}^{-}$or in $\Omega_{h}^{+}$;
(ii) For any edge $F, F$ has either vertices or the whole edge lying on $\Gamma$ if $F \cap \Gamma \neq \emptyset$.

We call $K$ an interface element if it intersects $\Gamma$ and denote the set of interface elements by $T_{h}^{*}$. For each $K \in T_{h}^{*}$, let $K^{-}=K \cap \Omega^{-}$and $K^{+}=K \cap \Omega^{+}$. Because $\Gamma$ is $C^{2}$ smooth, it implies either meas $\left(K^{-}\right) \leq c h_{K}^{3}$ or meas $\left(K^{+}\right) \leq c h_{K}^{3}$. Throughout this paper, we will use $\tilde{K}$ to denote one of the two subregions $K^{-}$and $K^{+}$which satisfies meas $\left(K^{s}\right) \leq$ $c h_{K}^{3}$. Here and later, $c$ denotes a generic positive constant independent of $h$ but may take different values at different occasions.


Fig. 2. An example of the triangulation
On triangulation $T_{h}$ we construct the piecewise $P_{1}$ linear conforming finite element space $V_{h}$ such that $V_{h} \subset H_{0}^{1}(\Omega) \cap$ $C(\bar{\Omega})$, and define $\Pi_{h}: H^{1}(\Omega) \rightarrow V_{h}$ to be the associated interpolation operator.

The corresponding discrete form of (1)-(2) reads as: find $\left(y_{h}, u_{h}\right) \in V_{h} \times U_{a d}$, such that

$$
\begin{equation*}
\min _{u_{h} \in U_{a d}} J_{h}\left(y_{h}, u_{h}\right)=\frac{1}{2}\left\|y_{h}-y_{d}\right\|_{0, \Omega}^{2}+\frac{\alpha}{2}\left\|u_{h}\right\|_{0, \Omega}^{2} \tag{8}
\end{equation*}
$$

subject to

$$
\begin{equation*}
a_{h}\left(y_{h}, v_{h}\right)=\left(u_{h}, v_{h}\right), \forall v_{h} \in V_{h}, \tag{9}
\end{equation*}
$$

where $a_{h}\left(y_{h}, v_{h}\right)=\sum_{K \in T_{h}} \int_{K} \beta_{h} \nabla y_{h} \nabla v_{h} d x, \beta_{h}=\beta^{s}$ if $K \subset \Omega_{h}^{s}$.

Similar to (5), we seek a unique solution $\left(y_{h}, p_{h}, u_{h}\right)$ satisfying the following discrete optimality conditions:

$$
\begin{cases}a_{h}\left(y_{h}, v_{h}\right)=\left(u_{h}, v_{h}\right), & \forall v_{h} \in V_{h}  \tag{10}\\ a_{h}\left(p_{h}, v_{h}\right)=\left(y_{h}-y_{d}, v_{h}\right), & \forall v_{h} \in V_{h} \\ \left(\alpha u_{h}+p_{h}, v_{h}-u_{h}\right) \geq 0, & \forall v_{h} \in U_{a d}\end{cases}
$$

where the optimal control $u_{h}$ will be solved from the adjoint state $p_{h}$,

$$
\begin{align*}
u_{h} & =P_{U_{a d}}\left\{-\frac{1}{\alpha} p_{h}\right\} \\
& =\min \left\{b(x), \max \left(a(x),-\frac{1}{\alpha} p_{h}\right)\right\} . \tag{11}
\end{align*}
$$

The following lemma has been presented in [5], which plays an important role in our theoretical analysis.

Lemma 2.1 Let $f \in L^{2}(\Omega)$, and $\Omega_{0} \in \Omega$ be a neighborhood of the interface $\Gamma$. Suppose that $\varphi \in Y \cap W^{1, \infty}\left(\Omega^{-} \cap \Omega_{0}\right) \cap$ $W^{1, \infty}\left(\Omega^{+} \cap \Omega_{0}\right)$ and $\varphi_{h} \in V_{h}$ are solutions of

$$
\begin{equation*}
a(\varphi, v)=(f, v), \quad \forall v \in H_{0}^{1}(\Omega) \tag{12}
\end{equation*}
$$

and

$$
\begin{equation*}
a_{h}\left(\varphi_{h}, v_{h}\right)=\left(f, v_{h}\right), \quad \forall v_{h} \in V_{h}, \tag{13}
\end{equation*}
$$

respectively. Then, there hold the following error estimate results:

$$
\begin{equation*}
\left|a\left(\varphi, v_{h}\right)-a_{h}\left(\varphi_{h}, v_{h}\right)\right| \leq c h\|\varphi\|_{Y, \Omega} \tag{14}
\end{equation*}
$$

and

$$
\begin{array}{ll}
\left\|\varphi-\Pi_{h} \varphi\right\|_{0, \Omega}+h\left|\varphi-\Pi_{h} \varphi\right|_{1, \Omega} & \leq c h^{2}\|\varphi\|_{Y, \Omega}, \\
\left\|\varphi-\varphi_{h}\right\|_{0, \Omega}+h\left|\varphi-\varphi_{h}\right|_{1, \Omega} & \leq c h^{2}\|\varphi\|_{Y, \Omega} \tag{15}
\end{array}
$$

where $\|\varphi\|_{Y, \Omega}:=\sqrt{\|\varphi\|_{1, \Omega}^{2}+|\varphi|_{2, \Omega^{+}}^{2}+|\varphi|_{2, \Omega^{-}}^{2}}$.

## III. OPTIMAL ORDER ERROR ESTIMATES

This section proceeds in two steps. First, we present optimal order error estimates and detailed proof of the state $y$ and adjoint state $p$ in $L^{2}$-norm. Second, the optimal order error estimate of the state $y$ and adjoint state $p$ in the brokenenergy norm will be proved in Theorem 3.2.

Theorem 3.1. Let $(u, y, p) \in U_{a d} \times Y \times Y$ and $\left(u_{h}, y_{h}, p_{h}\right) \in$ $U_{a d} \times V_{h} \times V_{h}$ be the solutions of (1) and (8), respectively. Then, there holds the following error estimate:

$$
\begin{equation*}
\left\|u-u_{h}\right\|_{0, \Omega}+\left\|y-y_{h}\right\|_{0, \Omega}+\left\|p-p_{h}\right\|_{0, \Omega} \leq c h^{2} . \tag{16}
\end{equation*}
$$

Proof. Replacing $v$ and $v_{h}$ with $u_{h}$ and $u$ in the inequalities of (5) and (10) yields

$$
\begin{equation*}
\left(\alpha u+p, u-u_{h}\right) \leq 0 \tag{17}
\end{equation*}
$$

and

$$
\begin{equation*}
\left(\alpha u_{h}+p_{h}, u_{h}-u\right) \leq 0 \tag{18}
\end{equation*}
$$

Then, it follows from summing up the above two inequalities and Lemma 2.1 that

$$
\begin{align*}
& \alpha\left\|u-u_{h}\right\|_{0, \Omega}^{2} \\
& \leq\left(u_{h}-u, p-p_{h}\right) \\
& =\left(u_{h}-u, p-p_{h}(y)\right)+\left(u_{h}-u, p_{h}(y)-p_{h}\right) \\
& =\left(u_{h}-u, p-p_{h}(y)\right)+a_{h}\left(y_{h}-y_{h}(u), p_{h}(y)-p_{h}\right) . \tag{19}
\end{align*}
$$

The first term on the right hand side of the above inequality can be estimated as follows:

$$
\begin{align*}
& \left(u_{h}-u, p-p_{h}(y)\right) \\
& \leq \frac{\alpha}{2}\left\|u_{h}-u\right\|_{0, \Omega}^{2}+\frac{1}{2 \alpha}\left\|p-p_{h}(y)\right\|_{0, \Omega}^{2} \tag{20}
\end{align*}
$$

Then, we are going to estimate the second term on the right hand side of (19). Actually, we have

$$
\begin{align*}
& a_{h}\left(y_{h}-y_{h}(u), p_{h}(y)-p_{h}\right) \\
& =a_{h}\left(y_{h}-y_{h}(u), p_{h}(y)\right)-a_{h}\left(y_{h}-y_{h}(u), p_{h}\right) \\
& =\left(y_{h}-y_{h}(u), y-y_{d}\right)-\left(y_{h}-y_{h}(u), y_{h}-y_{d}\right) \\
& =\left(y_{h}-y_{h}(u), y-y_{h}\right)  \tag{21}\\
& =\left(y_{h}-y, y-y_{h}\right)+\left(y-y_{h}(u), y-y_{h}\right) \\
& \leq \frac{1}{2}\left\|y-y_{h}(u)\right\|_{0, \Omega}^{2}-\frac{1}{2}\left\|y-y_{h}\right\|_{0, \Omega}^{2},
\end{align*}
$$

where $y_{h}(u) \in V_{h}$ and $p_{h}(y) \in V_{h}$ are the solutions of

$$
\begin{equation*}
a_{h}\left(y_{h}(u), v_{h}\right)=\left(u, v_{h}\right), \quad \forall v_{h} \in V_{h}, \tag{22}
\end{equation*}
$$

and

$$
\begin{equation*}
a_{h}\left(p_{h}(y), v_{h}\right)=\left(y-y_{d}, v_{h}\right), \quad \forall v_{h} \in V_{h}, \tag{23}
\end{equation*}
$$

respectively.

Summarizing the above two inequalities and substituting it into (19) lead to

$$
\begin{align*}
& \alpha\left\|u-u_{h}\right\|_{0, \Omega}^{2}+\left\|y-y_{h}\right\|_{0, \Omega}^{2} \\
& \leq\left\|p-p_{h}(y)\right\|_{0, \Omega}^{2}+\alpha\left\|y-y_{h}(u)\right\|_{0, \Omega}^{2} \tag{24}
\end{align*}
$$

Noticing that $p_{h}(y)$ and $y_{h}(u)$ are standard finite element approximations of $p$ and $y$. As a consequence, by Lemma 2.1, we have

$$
\begin{equation*}
\left\|p-p_{h}(y)\right\|_{0, \Omega} \leq c h^{2}\|p\|_{Y, \Omega} \tag{25}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\|y-y_{h}(u)\right\|_{0, \Omega} \leq c h^{2}\|y\|_{Y, \Omega} \tag{26}
\end{equation*}
$$

Combining (24), (25) and (26) gives that

$$
\begin{equation*}
\left\|u-u_{h}\right\|_{0, \Omega}+\left\|y-y_{h}\right\|_{0, \Omega} \leq c h^{2} . \tag{27}
\end{equation*}
$$

In the following, we consider the estimate of $\left\|p-p_{h}\right\|_{0, \Omega}$. By the definition of bilinear form $a_{h}(\cdot, \cdot)$ and (10), there exists a positive number $c_{0}$ such that

$$
\begin{align*}
c_{0}\left\|p_{h}(y)-p_{h}\right\|_{0, \Omega}^{2} & \leq a_{h}\left(p_{h}(y)-p_{h}, p_{h}(y)-p_{h}\right) \\
& =\left(p_{h}(y)-p_{h}, y-y_{h}\right) \\
& \leq c\left\|p_{h}(y)-p_{h}\right\|_{0, \Omega}\left\|y-y_{h}\right\|_{0} \\
& \leq c h^{2}\left\|p_{h}(y)-p_{h}\right\|_{0, \Omega} \tag{28}
\end{align*}
$$

which implies that

$$
\begin{equation*}
\left\|p_{h}(y)-p_{h}\right\|_{0, \Omega} \leq c h^{2} . \tag{29}
\end{equation*}
$$

Combining (25) and (29) yields

$$
\begin{equation*}
\left\|p-p_{h}\right\|_{0, \Omega} \leq c h^{2} \tag{30}
\end{equation*}
$$

The proof is completed.
Now we are ready to derive the optimal order error estimates for the state $y$ and adjoint state $p$ in the broken energy norm.

Theorem 3.2. Under the assumption of Theorem 3.1, there hold the following optimal order error estimates for state $y$ and adjoint state $p$ :

$$
\begin{equation*}
\left|y-y_{h}\right|_{1, \Omega} \leq c h \tag{31}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|p-p_{h}\right|_{1, \Omega} \leq c h \tag{32}
\end{equation*}
$$

respectively.
Proof. First of all, let $\beta_{*}=\min \left\{\beta^{-}, \beta^{+}\right\}$. We have

$$
\begin{align*}
& \beta_{*}\left|y-y_{h}\right|_{1, \Omega}^{2} \leq a_{h}\left(y-y_{h}, y-y_{h}\right) \\
& =a_{h}\left(y-y_{h}, y-\Pi_{h} y\right)+a_{h}\left(y-y_{h}, \Pi_{h} y-y_{h}\right) \tag{33}
\end{align*}
$$

The bound of the first term of (33) can be found directly from Schwarz inequality and the standard approximation theory, i.e.,

$$
\begin{align*}
a_{h}\left(y-y_{h}, y-\Pi_{h} y\right) & \leq c\left|y-y_{h}\right|_{1, \Omega}\left|y-\Pi_{h} y\right|_{1, \Omega} \\
& \leq \frac{\beta_{*}}{4}\left|y-y_{h}\right|_{1, \Omega}^{2}+c\left|y-\Pi_{h} y\right|_{1, \Omega}^{2} \\
& \leq c h^{2}+\frac{\beta_{*}}{4}\left|y-y_{h}\right|_{1, \Omega}^{2} . \tag{34}
\end{align*}
$$

The estimation of the second term of (33) follows from the results of Lemma 2.1, Theorem 3.1, and the standard approximation theory

$$
\begin{align*}
& a_{h}\left(y-y_{h}, \Pi_{h} y-y_{h}\right) \\
& \leq\left(u-u_{h}, \Pi_{h} y-y_{h}\right)+c h\|y\|_{Y, \Omega}\left|\Pi_{h} y-y_{h}\right|_{1, \Omega} \\
& \leq\left\|u-u_{h}\right\|_{0, \Omega}\left\|\Pi_{h} y-y_{h}\right\|_{0, \Omega}+c h\|y\|_{Y, \Omega}\left|\Pi_{h} y-y_{h}\right|_{1, \Omega} \\
& \leq c\left(h^{4}+h^{2}\|y\|_{Y, \Omega}^{2}+\left\|\Pi_{h} y-y\right\|_{1, \Omega}^{2}\right)+\frac{\beta_{*}}{4}\left|y-y_{h}\right|_{1, \Omega}^{2} \\
& \leq c h^{2}+\frac{\beta_{*}}{4}\left|y-y_{h}\right|_{1, \Omega}^{2} . \tag{35}
\end{align*}
$$

Summarize the above two inequalities into (33) yields (31). Similarly, for the adjoint state $p$, using again Lemma 2.1 and Theorem 3.1, we have

$$
\begin{align*}
& \beta_{*}\left|p-p_{h}\right|_{1, \Omega}^{2} \\
& \leq a_{h}\left(p-p_{h}, p-p_{h}\right) \\
& =a_{h}\left(p-p_{h}, p-\Pi_{h} p\right)+a_{h}\left(p-p_{h}, \Pi_{h} p-p_{h}\right) \\
& \leq a_{h}\left(p-p_{h}, p-\Pi_{h} p\right)+\left(y-y_{h}, \Pi_{h} p-p_{h}\right) \\
& +c h\left|\Pi_{h} p-p_{h}\right|_{1, \Omega} \\
& \leq c\left|p-p_{h}\right|_{1, \Omega}\left|p-\Pi_{h} p\right|_{1, \Omega}+\left\|y-y_{h}\right\|_{0, \Omega}\left|\Pi_{h} p-p_{h}\right|_{1, \Omega} \\
& +c h\left|\Pi_{h} p-p_{h}\right|_{1, \Omega} \\
& \leq c h\left|p-p_{h}\right|_{1, \Omega}+c h\left(\left|\Pi_{h} p-p\right|_{1, \Omega}+\left|p-p_{h}\right|_{1, \Omega}\right) \\
& \leq c h\left|p-p_{h}\right|_{1, \Omega}+c h^{2}\|p\|_{Y, \Omega} \\
& \leq c h^{2}+\frac{\beta_{*}}{2}\left|p-p_{h}\right|_{1, \Omega}^{2}, \tag{36}
\end{align*}
$$

which gives (32). The proof is thus completed.

## IV. Numerical experiment

This section will provide some numerical results for the elliptic interface control problem to verify the correctness of the theorems given in the previous section.

In this example we choose $\alpha=1$ and the computation domain as $\Omega=[-1,1] \times[-1,1]$, the interface $\Gamma$ is a circle centered at the origin with radius being $r_{0}=0.5 . \Omega^{-}=$ $\left\{\left(x_{1}, x_{2}\right) \mid x_{1}^{2}+x_{2}^{2} \leq 0.5\right\}, \Omega^{+}=\Omega-\Omega^{-}$.

The admissible control set $U_{a d}$ is given as

$$
\begin{equation*}
U_{a d}=\{v \in U:-1 \leq v \leq 1, \text { a.e. in } \Omega\} . \tag{37}
\end{equation*}
$$

We take the optimal state and adjoint state as

$$
y=\left\{\begin{array}{l}
u^{-}=\frac{\left(x_{1}^{2}+x_{2}^{2}\right)^{3 / 2}}{\beta^{-}}, \quad \text { in } \Omega^{-},  \tag{38}\\
u^{+}=\frac{\left(x_{1}^{2}+x_{2}^{2}\right)^{3 / 2}}{\beta^{-}}+\left(\frac{1}{\beta^{-}}-\frac{1}{\beta^{+}}\right) r_{0}^{3}, \quad \text { in } \Omega^{+},
\end{array}\right.
$$

and
$p= \begin{cases}p^{-}=\frac{5\left(x_{1}^{2}+x_{2}^{2}-r_{0}^{2}\right)\left(1-x_{1}\right)\left(x_{1}+1\right)\left(x_{2}-1\right)\left(x_{2}+1\right)}{\beta^{-}}, & \text {in } \Omega^{-}, \\ p^{+}=\frac{5\left(x_{1}^{2}+x_{2}^{2}-r_{0}^{2}\right)\left(1-x_{1}\right)\left(x_{1}+1\right)\left(x_{2}-1\right)\left(x_{2}+1\right)}{\beta^{+}}, & \text {in } \Omega^{+},\end{cases}$
respectively.
The optimal control could be expressed as

$$
\begin{align*}
u(x) & =P_{U_{a d}}\{-p(x)\} \\
& =\min \{1, \max (-1,-p(x))\} \tag{40}
\end{align*}
$$

Then the functions $f$ and $y_{d}$ can be determined the above functions accordingly.
In this experiment, we fix $\beta^{-}=-1$, and consider $\beta^{+}=5$ and $\beta^{+}=50$ as two cases. We first approximate the circle $\Gamma$ by a polygon, and then give triangular subdivision to these
two domains separately. A uniform triangle grid mesh is thus completed. The error estimates and convergence orders of the control, state and adjoint state are shown in the following Tables 1-4 for $\beta^{+}=5$ and $\beta^{+}=50$, and the convergence rates are reported in Figures 3-6, where $N$ denotes the number of the elements, "order" represents the convergence order which is evaluated by

$$
\begin{equation*}
\text { Order }=\frac{1}{\log \left(N_{2} / N_{1}\right)^{1 / 2}} \log \frac{\left\|u-u_{N_{1}}\right\|_{i, \Omega}}{\left\|u-u_{N_{2}}\right\|_{i, \Omega}}, \tag{41}
\end{equation*}
$$

here, $\left\|u-u_{N}\right\|_{i}$ is a special norm for $i=0,1$.

| Table 1 The errors and convergence orders in $L^{2}$-norm with $\beta^{+}=5$ |  |  |  |
| :---: | :---: | :---: | :---: |
| $N$ | $\left\\|u-u_{h}\right\\|_{0, \Omega}$ | $\left\\|y-y_{h}\right\\|_{0, \Omega}$ | $\left\\|p-p_{h}\right\\|_{0, \Omega}$ |
| 14 | 0.318638863 | 1.189625451 | 0.097867163 |
| order | $/$ | $/$ | $/$ |
| 72 | 0.065900590 | 0.248701791 | 0.021369779 |
| order | 1.92465 | 1.91149 | 1.85836 |
| 322 | 0.013276842 | 0.061789328 | 0.005924887 |
| order | 2.13918 | 1.85932 | 1.71284 |
| 1458 | 0.002700346 | 0.012417775 | 0.001352403 |
| order | 2.10908 | 2.12492 | 1.95631 |
| 2982 | 0.001287701 | 0.005823426 | 0.000638279 |
| order | 2.06986 | 2.11659 | 2.09876 |


| Table 2 The errors and convergence orders in energy norm with $\beta^{+}=5$ |  |  |
| :---: | :---: | :---: |
| $N$ | $\left\|y-y_{h}\right\|_{1, \Omega}$ | $\left\|p-p_{h}\right\|_{1, \Omega}$ |
| 14 | 0.164654461 | 0.003396546 |
| order | $/$ | $/$ |
| 72 | 0.073176671 | 0.001588027 |
| order | 0.99043 | 0.92851 |
| 322 | 0.032943245 | 0.000887296 |
| order | 1.06562 | 0.77719 |
| 1458 | 0.018774017 | 0.000390204 |
| order | 0.74465 | 1.08790 |
| 2982 | 0.012140199 | 0.000271377 |
| order | 1.21854 | 1.01508 |



Fig. 3. Convergence rates of $L^{2}$ norm with $\beta^{+}=5$


Fig. 4. Convergence rates of energy norm with $\beta^{+}=5$

| Table 3 The errors and convergence orders in $L^{2}$-norm with $\beta^{+}=50$ |  |  |  |
| :---: | :---: | :---: | :---: |
| $N$ | $\left\\|u-u_{h}\right\\|_{0, \Omega}$ | $\left\\|y-y_{h}\right\\|_{0, \Omega}$ | $\left\\|p-p_{h}\right\\|_{0, \Omega}$ |
| 14 | 0.678309069 | 3.631292383 | 0.171599428 |
| order | $/$ | $/$ | $/$ |
| 72 | 0.146790976 | 0.81557836 | 0.039212493 |
| order | 1.86930 | 1.82394 | 1.80283 |
| 322 | 0.039629437 | 0.214722199 | 0.010928453 |
| order | 1.74838 | 1.78192 | 1.70591 |
| 1458 | 0.00858952 | 0.049549885 | 0.003454323 |
| order | 2.02484 | 1.94186 | 1.52522 |
| 2982 | 0.004611587 | 0.023641828 | 0.001747589 |
| order | 1.73849 | 2.06830 | 1.90458 |


| Table 4 The errors and convergence orders in energy norm with $\beta^{+}=50$ |  |  |
| :---: | :---: | :---: |
| $N$ | $\left\|y-y_{h}\right\|_{1, \Omega}$ | $\left\|p-p_{h}\right\|_{1, \Omega}$ |
| 14 | 0.537165421 | 0.062955460 |
| order | $/$ | $/$ |
| 72 | 0.227192916 | 0.024646492 |
| order | 1.05093 | 1.14532 |
| 322 | 0.132958099 | 0.013176262 |
| order | 0.71536 | 0.83614 |
| 1458 | 0.066483736 | 0.005838989 |
| order | 0.91782 | 1.07777 |
| 2982 | 0.043656470 | 0.003771959 |
| order | 1.17565 | 1.22137 |
|  |  |  |



Fig. 5. Convergence rates of $L^{2}$ norm with $\beta^{+}=50$


Fig. 6. Convergence rates of energy norm with $\beta^{+}=50$

We can see that this linear body-fitted method could achieve a optimal order convergence, which almost coincides with our theoretical analysis.
Remark 1. It is worth mentioning that the fact $\left.\left(\nabla v_{h}\right)\right|_{K}=$ constant for all $v_{h} \in V_{h}$ is crucial in the error analysis, which implies that the idea in the error analysis can be apply to $P_{1}$ - nonconforming triangular element [9] and $P_{1}$-rectangular element [22]. However, this approach could not be generalized to higher order elements.

Remark 2. With properly handling the time variable, the results obtained in this paper could be extended to some timedependent OCPs, such as parabolic and hyperbolic interface control problems.

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